Robust time-domain source stepping for DC-solution of circuit equations

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Summary. Most analyses of circuit equations start with solving the steady-state (DC) solution. In several cases this can be very hard. We present a novel time domain source stepping procedure to obtain a DC solution of circuit equations. The source stepping procedure is automatically adaptive. Controlled sources can be elegantly dealt with. The method can easily be combined with existing pseudo-transient procedures. The method is robust and efficient.

1 Introduction

The circuit equations can be written as [5, 10]

$$\frac{\mathrm{d}}{\mathrm{d}t}\mathbf{q}(\mathbf{x}) + \mathbf{j}(\mathbf{x}) + \mathbf{s}(t, \mathbf{x}) = 0 \tag{1}$$

Here $\mathbf{s}(t, \mathbf{x})$ represents the specifications of the sources. The unknown $\mathbf{x} = \mathbf{x}(t)$ consists of nodal voltages and of currents through voltage defined elements. We assume that $\mathbf{q}(0) = 0$, and $\mathbf{j}(0) = 0$.

The steady state solution, which is called DC-solution (Direct Current solution), x_{DC} , satisfies

$$\mathbf{j}(\mathbf{x}_{\mathrm{DC}}) + \mathbf{s}(0, \mathbf{x}_{\mathrm{DC}}) = 0.$$
 (2)

Usually, and already hinted by setting t = 0 in (2), the DC-solution provides the initial value for the transient problem (1). In general, the problem (2) is non-linear. How to solve this problem is the subject of this note. The importance of the DC-problem lies in the fact that the DC-solution is crucial as starting solution for a number of next analyses (transient analysis, AC analysis, Harmonic Balance analysis, Periodic Steady-State analysis). In general, (1) forms a system of Differential-Algebraic Equations (DAEs). With $\mathscr{C} =$ $\frac{\partial q(x)}{\partial x}$ $\left\| \sum_{\mathbf{x}=\mathbf{x}_{DC}} \text{ and } \mathscr{G} = \frac{\partial \mathbf{j}(\mathbf{x})}{\partial \mathbf{x}} \right\|_{\mathbf{x}=\mathbf{x}_{DC}}. \text{ We assume that } \lambda \mathscr{C} +$ \mathscr{G} is non-singular for λ in some neighbourhood of 0 (may be excluding $\lambda = 0$). To solve the equations Newton's method, or variants, may be applied [3,5,8], which can be combined with g_{\min} -stepping, in which linear conductors g are placed parallel to the nonlinear part inside each transistor (device). Iteratively $g \downarrow g_{\min}$, after which the Newton counter is increased.

Another approach is Pseudo-Transient [2]. In Pseudo-Transient (PT) one can use relaxed tolerances for the Newton process and for the time step control procedure. Also this can be combined with g_{min} -stepping during each time step. In PT one has to provide a non-trivial initial solution. A new procedure is decribed in the next section. Other methods are: temperature stepping, source stepping (the sources are iteratively increased to their final value), homotopy methods, or optimization [1,4,7,9–12].

2 Time-domain Source Stepping

Usually, in Source Stepping one introduces a parameter λ and considers the problem

$$\mathbf{j}(\mathbf{x}(\boldsymbol{\lambda})) + \boldsymbol{\lambda}\mathbf{s}(0, \mathbf{x}(\boldsymbol{\lambda})) = 0.$$
(3)

In this case it is assumed that for $\lambda = 0$ the problem (3) is easily solved so that in the end the original problem is solved. The same parameter λ is applied to all sources *s* in the circuit. In general, for each value of λ a nonlinear problem has to be solved.

We introduce a time-domain variant (SSPT) that offers an automatic continuation process, based on PT and adapting the transient stepsize and the λ stepsize at the same time.

We define a time t = T at which we want to have solved the original DC-problem. We also introduce a time $T_{\alpha} = \alpha T$ (by default $\alpha = 0.5$) at which ordinary PT will start simulation using the sources as in the original DC-problem, i.e. using $\lambda = 1$ and where PT integrates from T_{α} to T', where $T' \leq T$ is the point where all transient effects have become negligible (see also Fig. 1).

On the interval $[0, T_{\alpha}]$, a special PT integration is performed with the function $\lambda(t) = t/T_{\alpha}$. Hence, at each time step, also the actual applied source values change. The interval $[0, T_{\alpha}]$ is the switch-on interval, the interval $[T_{\alpha}, T]$ is the interval to damp-out transient effects. On both intervals PT uses an automatic time step determination procedure. On the interval $[T_{\alpha}, T]$ an ordinary PT procedure is executed. Hence, if, at some time point, the Newton iterative process does not converge, a re-integration will be done with a smaller stepsize. Recursion in controlled sources asks

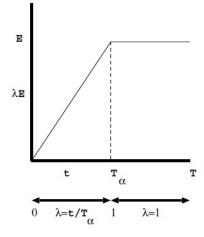


Fig. 1. On $[0, T_{\alpha}]$ a time-dependent voltage source $\lambda(t)\mathbf{E}$ is used where $\lambda(t) = t/T_{\alpha}$. On $[T_{\alpha}, T]$ we have $\lambda \equiv 1$.

for a modification in (3). An expression for a controlled voltage source $E_1(0, 1)$ may look like

$$V(E_1) = 5 + 4I(E_1) + [6V(R_1) + 7I(E_2) + 12]^2$$
(4)

It is controlled by the controlling "ev's" (electrical variables) $I(E_1)$, $V(R_1)$, and $I(E_2)$. We write the expression for the applied value $V(E_1)$ as

$$V(E_1) = \Psi(\text{ev}_1, \text{ev}_2, \dots, \text{ev}_n) \tag{5}$$

As value during the source stepping at time t on $[0, T_{\alpha}]$ we propose to take

$$V(E_1) = \tilde{\psi}(\text{ev}_1, \dots, \text{ev}_n), \text{ where } (6)$$

$$\tilde{\psi}(\text{ev}_1, \dots, \text{ev}_n) = \psi(\text{ev}_1, \dots, \text{ev}_n) + (\lambda(t) - 1)\psi(0, \dots, 0). \quad (7)$$

Note that in (4), $\psi(0, ..., 0) = 149$. This value has to be calculated once. When in (4) E_2 is a controlled voltage source too, contributions to the Jacobian matrix are calculated by $\frac{\partial \Psi}{\partial \mathbf{x}} = \frac{\partial \Psi}{\partial ev_i} \frac{\partial ev_i}{\partial \mathbf{x}}$, which gives recursion. Note that λ does not occur in the matrix. Clearly, for $\lambda = 0$ the applied voltage is zero (assuming starting from the zero solution, which implies that all ev's are zero), which makes the zero solution the exact solution. When $\lambda = 1$ the original voltage expression is used. Since our equations (1) are DAEs we remark that for all *t* the generated solution is consistent for the problem at hand. Because of the switchon and the damp-out phase the process mimics a real physical process.

3 Results

We tested the SSPT on a set of difficult problems where parameters were swept (temperature, and statistics). The SSPT was always convergent (without needing g_{\min} -iteration). It was 1-13 times faster than Newton-Raphson (that sometimes needed internal g_{\min} iteration). Normal PT was less robust than SSPT. Further improvements in the time-domain integrations, after starting with a proper \mathbf{x}_{DC} , have been tuned to fault analysis [6].

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